Acoustic Models: Analytic Solution Techniques

"He suffered from paralysis by analysis," Harold S. Geneen

Linear Wave Equations for Sound

Pressure:

$$\frac{\partial^2 \hat{p}}{\partial t^2} = c^2 \Delta \hat{p}$$

Boundary Conditions

Initial Conditions

Velocity:

$$\frac{\partial^2 \hat{u}}{\partial t^2} = c^2 \Delta \hat{u}$$

Boundary Conditions

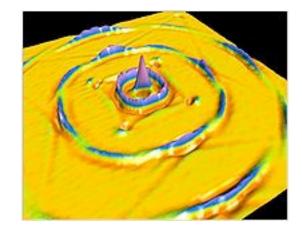
Initial Conditions

Potential:

$$\frac{\partial^2 \phi}{\partial t^2} = c^2 \Delta \phi$$

Boundary Conditions

Initial Conditions

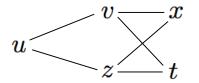


Model:

$$\begin{split} &\frac{\partial^2 u}{\partial t^2} = c^2 \frac{\partial^2 u}{\partial x^2} & -\infty < x < \infty \ , \ t > 0 \\ &u(0,x) = f(x) \ , \ u_t(0,x) = g(x) & -\infty < x < \infty \end{split}$$

New Independent Variables:

v = x + ct, z = x - ct



Thus

$$u_x = u_v v_x + u_z z_x = u_v + u_z$$

and

$$u_{xx} = (u_v + u_z)_x = (u_v + u_z)_v v_x + (u_v + u_z)_z z_x = u_{vv} + 2u_{vz} + u_{zz}$$

Similarly,

$$u_{tt} = c^2(u_{vv} - 2u_{vz} + u_{zz})$$

Transformed Differential Equation:

$$u_{vz} \equiv \frac{\partial^2 u}{\partial v \partial z} = 0$$

$$\Rightarrow \frac{\partial u}{\partial v} = h(v)$$

$$\Rightarrow u(v, z) = \int_0^v h(s) ds + B(z)$$

$$\Rightarrow u(v, z) = A(v) + B(z)$$

where A(v) and B(z) are arbitrary functions that satisfy initial conditions.

General Solution:

$$u(t,x) = A(x+ct) + B(x-ct)$$

Initial Conditions:

$$u(0,x) = A(x) + B(x) = f(x)$$
, $u_t(0,x) = cA'(x) - cB'(x) = g(x)$

Second equation yields

$$A(x) - B(x) = \frac{1}{c} \int_0^x g(s) ds + D$$

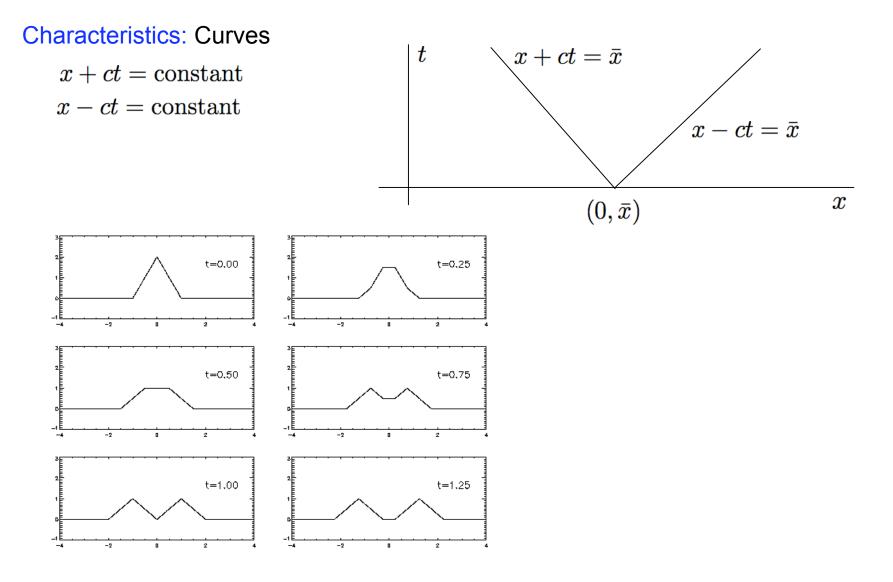
$$\Rightarrow \quad A(x) = \frac{1}{2} f(x) + \frac{1}{2c} \int_0^x g(s) ds + \frac{D}{2}$$

$$B(x) = \frac{1}{2} f(x) - \frac{1}{2c} \int_0^x g(s) ds - \frac{D}{2}$$

D'Alembert's Solution:

$$u(t,x) = rac{1}{2} \left[f(x+ct) + f(x-ct)
ight] + rac{1}{2c} \int_{x-ct}^{x+ct} g(s) ds$$

Physical Interpretation: Consider an infinite string with local initial conditions



Properties of Fourier Series

Fourier Series: Consider the representation

$$f(x) = \frac{a_0}{2} + \sum_{n=1}^{\infty} \left[a_n \cos\left(\frac{n\pi x}{L}\right) + b_n \sin\left(\frac{n\pi x}{L}\right) \right]$$

for a 2L periodic, continuous function f(x). If the series converges uniformly, it is the Fourier series for f(x).

Note:

$$\int_{-L}^{L} f(x)dx = a_0 L$$
$$\int_{-L}^{L} \sin\left(\frac{m\pi x}{L}\right) \sin\left(\frac{n\pi x}{L}\right) dx = \int_{-L}^{L} \cos\left(\frac{m\pi x}{L}\right) \cos\left(\frac{n\pi x}{L}\right) dx = 0 \quad \text{for } m \neq n$$
$$\int_{-L}^{L} \sin\left(\frac{m\pi x}{L}\right) \cos\left(\frac{n\pi x}{L}\right) dx = 0 \quad \text{for all } m, n$$

$$\int_{-L}^{L} \sin^2\left(\frac{m\pi x}{L}\right) dx = \cos^2\left(\frac{n\pi x}{L}\right) dx = L \quad \text{for } n \ge 1$$

Properties of Fourier Series

Fourier Coefficients:

$$a_n = \frac{1}{L} \int_{-L}^{L} f(x) \cos\left(\frac{n\pi x}{L}\right) dx \quad \text{for } n = 0, 1, 2, \dots$$
$$b_n = \frac{1}{L} \int_{-L}^{L} f(x) \sin\left(\frac{n\pi x}{L}\right) dx \quad \text{for } n = 1, 2, \dots$$

Theorem: Suppose that f and f' are piecewise continuous on the interval $-L \le x \le L$. Further, suppose that f is defined outside the interval $-L \le x \le L$ so that it is periodic with period 2L. Then f has a Fourier series

$$f(x) = \frac{a_0}{2} + \sum_{n=1}^{\infty} \left[a_n \cos\left(\frac{n\pi x}{L}\right) + b_n \sin\left(\frac{n\pi x}{L}\right) \right]$$

with the above coefficients. The Fourier series converges to f(x) at all points where f is continuous, and to [f(x+) + f(x-)]/2 at all points where f is discontinuous (see Boyce and DiPrima).

Separation of Variables

Model:

$$u_{tt} = c^2 u_{xx}$$
 , $0 < x < L$, $t > 0$
 $u(0,x) = f(x)$, $u_t(0,x) = g(x)$, $0 \le x \le L$
 $u(t,0) = u(t,L) = 0$, $t \ge 0$

Separation of Variables: Consider solutions of the form

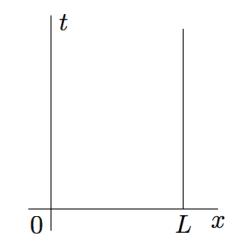
$$u(t, x) = X(x)T(t)$$

$$\Rightarrow X(x)\ddot{T}(t) = c^2 X''(x)T(t)$$

$$\Rightarrow \frac{X''(x)}{X(x)} = \frac{\ddot{T}(t)}{c^2 T(t)} = K$$

Consider first

$$X''(x) - KX(x) = 0$$
$$X(0) = X(L) = 0$$



Separation of Variables

Note:

$$\int_{0}^{L} \left[XX'' - KX^{2} \right] dx = \int_{0}^{L} \left[-(X')^{2} - KX^{2} \right] dx = 0$$

$$\Rightarrow \int_{0}^{L} [X'(x)^{2} + KX^{2}(x)] dx = 0$$

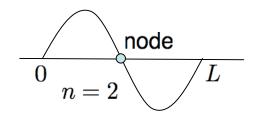
If $K \ge 0$, this implies X(x) = K = 0. Thus K < 0 so take $K = -\lambda^2, \lambda > 0$.

Boundary Value Problem: Helmholtz Equation

$$X''(x) + \lambda^2 X(x) = 0$$
$$X(0) = X(L) = 0$$

Solution: $X(x) = A\cos(\lambda x) + B\sin(\lambda x)$ $X(0) = 0 \Rightarrow A = 0$ $X(L) = 0 \Rightarrow \lambda L = n\pi$

Thus $X_n(x) = B_n \sin(\lambda_n x)$, $\lambda_n = \frac{n\pi}{L}$, $B_n \neq 0$ Modes or Eigenfunctions



Separation of Variables

Initial Value Problem Solution:

$$T_n(t) = C_n \cos(\lambda_n ct) + D_n \sin(\lambda_n ct)$$

General Solution to PDE:

$$u(t,x) = \sum_{n=1}^{\infty} [a_n \cos(\lambda_n ct) + b_n \sin(\lambda_n ct)] \sin(\lambda_n x)$$

Initial Conditions:

$$f(x) = \sum_{n=1}^{\infty} a_n \sin\left(\frac{n\pi x}{L}\right) \quad , \quad g(x) = \sum_{n=1}^{\infty} b_n \lambda_n c \sin\left(\frac{n\pi x}{L}\right)$$
$$\Rightarrow \quad a_n = \frac{2}{L} \int_0^L f(x) \sin\left(\frac{n\pi x}{L}\right) dx$$
$$b_n = \frac{2}{\lambda_n cL} \int_0^L g(x) \sin\left(\frac{n\pi x}{L}\right) dx$$

Example

Model: Take g(x) = 0 so $b_n = 0$

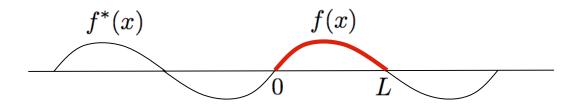
Note:

$$\cos(\lambda_n ct)\sin(\lambda_n x) = \frac{1}{2}[\sin(\lambda_n (x - ct)) + \sin(\lambda_n (x + ct))]$$

Solution:

$$u(t,x) = \frac{1}{2} \sum_{n=1}^{\infty} a_n \sin(\lambda_n(x-ct)) + \frac{1}{2} \sum_{n=1}^{\infty} a_n \sin(\lambda_n(x+ct))$$
$$\Rightarrow u(t,x) = \frac{1}{2} [f^*(x-ct) + f^*(x+ct)]$$

where f^* is the odd periodic extension of f with period 2L.



Relationship Between Wavelength and Frequency

Note:

- Vibrate structure with frequency f for time t
- Generate N = ft waves
- First wave travels distance ct
- Wavelength is ratio of distance to number of waves in this distance

Relationship:

$$\lambda = \frac{ct}{N} = \frac{ct}{ft}$$
$$\Rightarrow \lambda = \frac{c}{f}$$

